



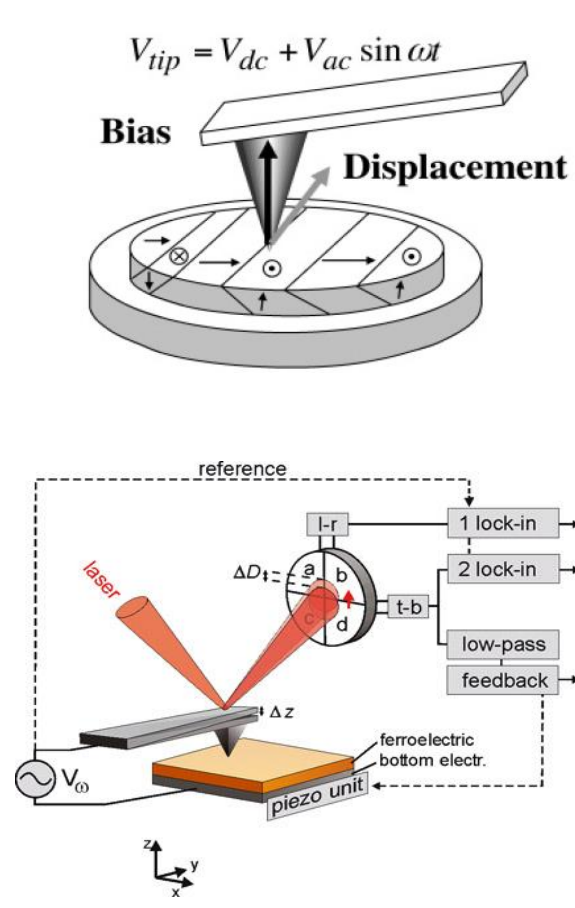
Ferroelastic Switching in PZT Thin Films: Phase-field Modeling

Zijian Hong, Jason Britson, Jia-Mian Hu, Long-Qing Chen

Department of Materials Science and Engineering, The Pennsylvania State University,
University Park, Pennsylvania 16802, USA

Background

Piezoresponse force microscopy (PFM) [1]



❖ Basic principle: Converse piezoelectric effect

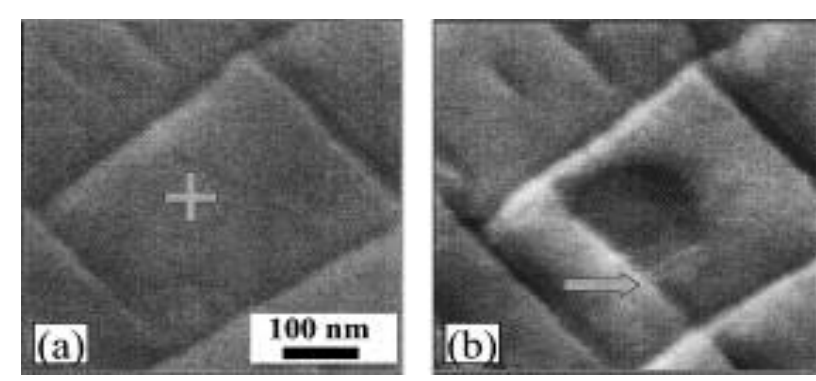
$$X_i = d_{ki} E_k$$

❖ Applications: domain imaging, domain patterning, studies of domain dynamics, spectroscopy and switching spectroscopy mapping

❖ Tip geometry: Lorentz distribution of potential

$$\varphi(x, y) = \varphi_0 \frac{\gamma^2}{\gamma^2 + (x - x_0)^2 + (y - y_0)^2}$$

PFM tip induced 90° switching [2]



❖ Local 90° switching has been observed in experiments

❖ Lack of detailed explanation and theoretical simulations

Methodology

Time-dependent Ginzburg-Landau equations [3]

$$\frac{\partial P_i(\mathbf{x}, t)}{\partial t} = -L \frac{\delta F}{\delta P_i(\mathbf{x}, t)}, i = 1, 2, 3$$

$$F = \int (f_{bulk} + f_{grad} + f_{elas} + f_{elec}) dV$$

Semi-implicit Fourier scheme [4]

$$\overline{P_i}^{n+1}(\mathbf{k}, t) = \frac{\overline{P_i}^n(\mathbf{k}, t) + \Delta t f_i^n(\mathbf{k}, t)}{1 + \Delta t k_{\omega}^2}$$

$$f_{bulk} = \frac{1}{2} \alpha_{ij} P_i P_j + \frac{1}{4} \gamma_{ijkl} P_i P_j P_k P_l + \dots$$

$$f_{grad} = \frac{1}{2} G_{ijkl} P_{i,j} P_{k,l}$$

$$f_{elas} = \frac{1}{2} c_{ijkl} (\varepsilon_{ij} - \varepsilon_{ij}^0) (\varepsilon_{kl} - \varepsilon_{kl}^0)$$

$$f_{elec} = -\frac{1}{2} \varepsilon_0 \kappa_{ij} E_i E_j - E_i P_i$$

Boundary conditions

$$\varphi_{bottom} = 0, \varphi_{top} = \varphi(x, y)$$

$$\sigma_{i3}|_{z=0} = 0$$

$$u_i|_{substrate/bottom} = 0$$

Additional conditions

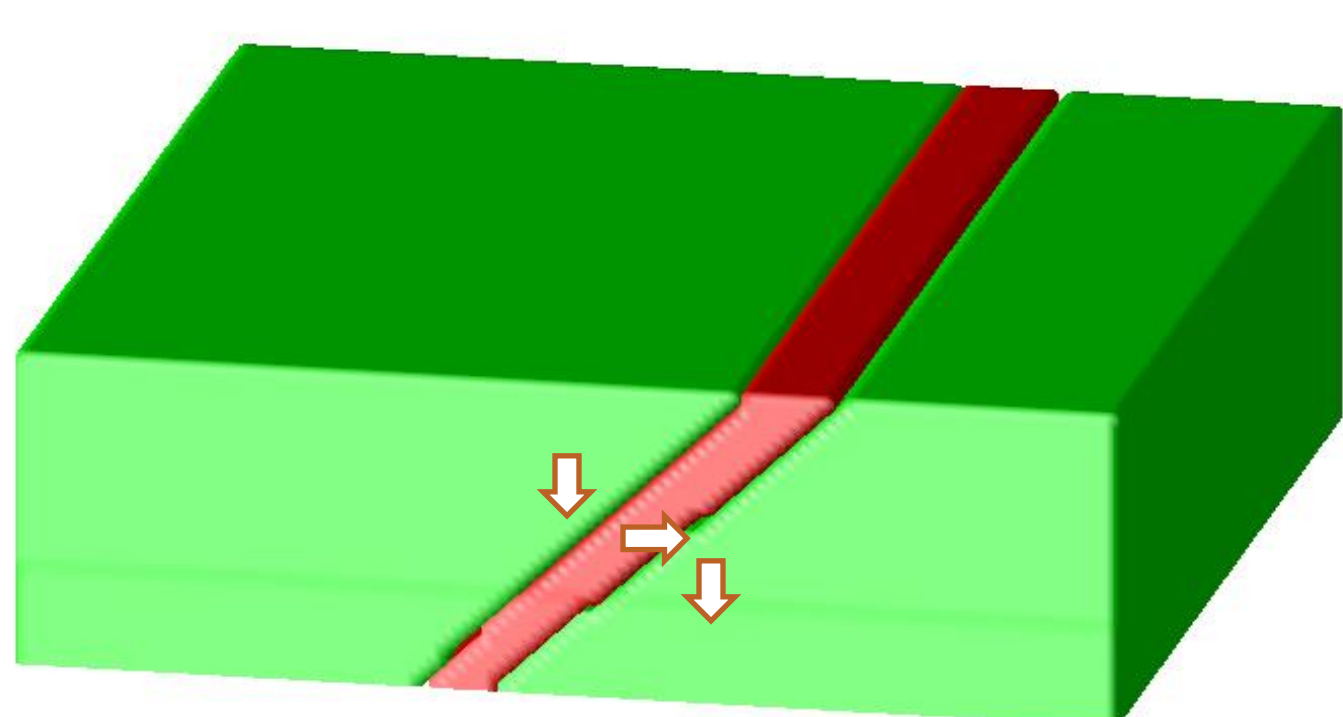
$$\varepsilon_0 \kappa_{ij} \frac{\partial^2 \varphi}{\partial x_i \partial x_j} = \frac{\partial P_i}{\partial x_j}$$

$$\sigma_{ij,j} = 0$$

$$\sigma_{ij} = c_{ijkl} (\varepsilon_{kl} - \varepsilon_{kl}^0)$$

Ferroelastic switching and domain wall influence

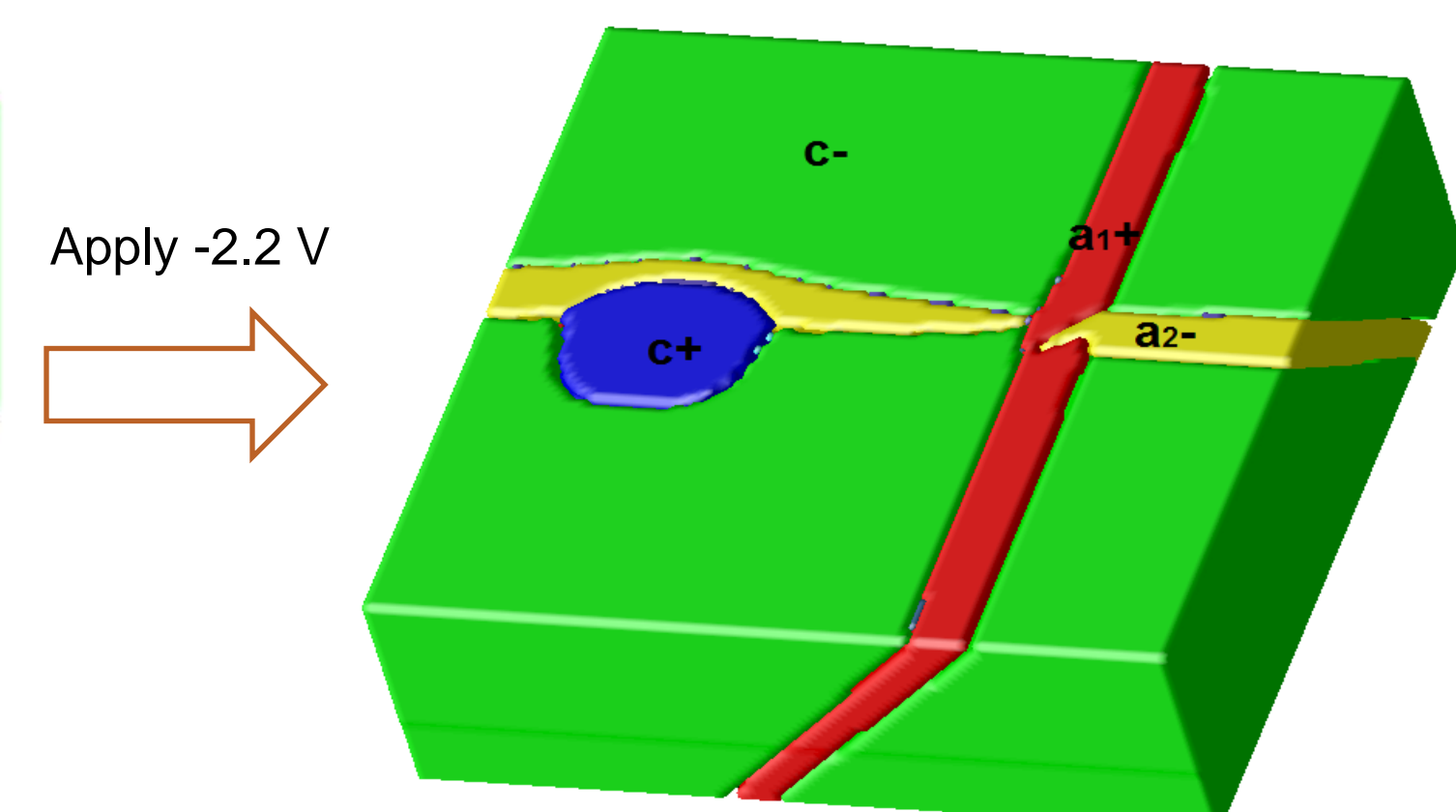
(1) Initial domain configuration



❖ We stabilize the a/c configuration as our initial domain

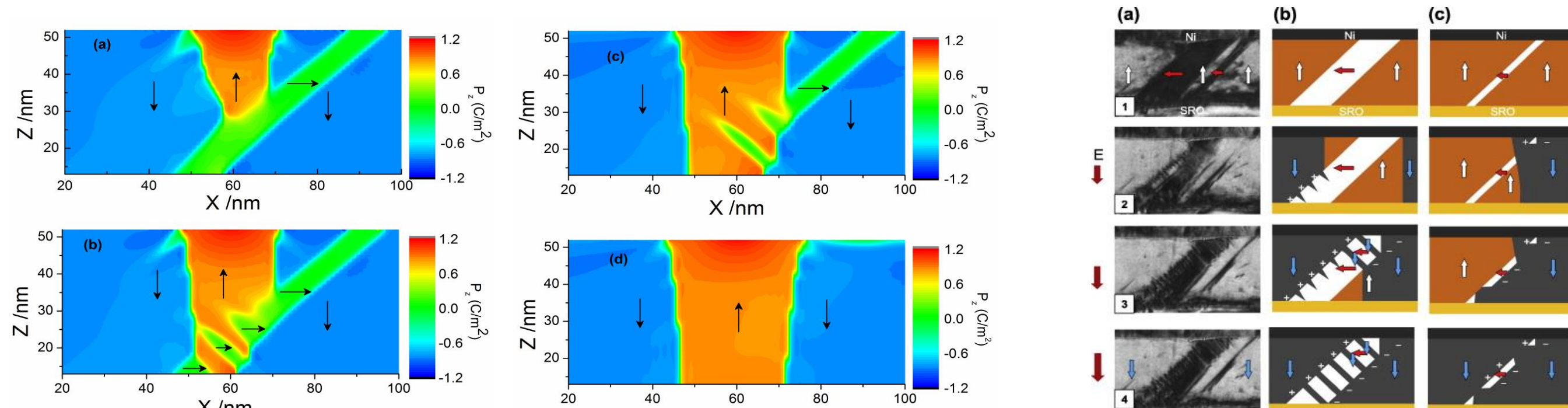
❖ The a domain tilts 45 degree with the c domain

(2) 90° switching far from domain wall



❖ When switched at high bias, 90° switching is observed along with local 180° switching

(3) 90° domain wall influence



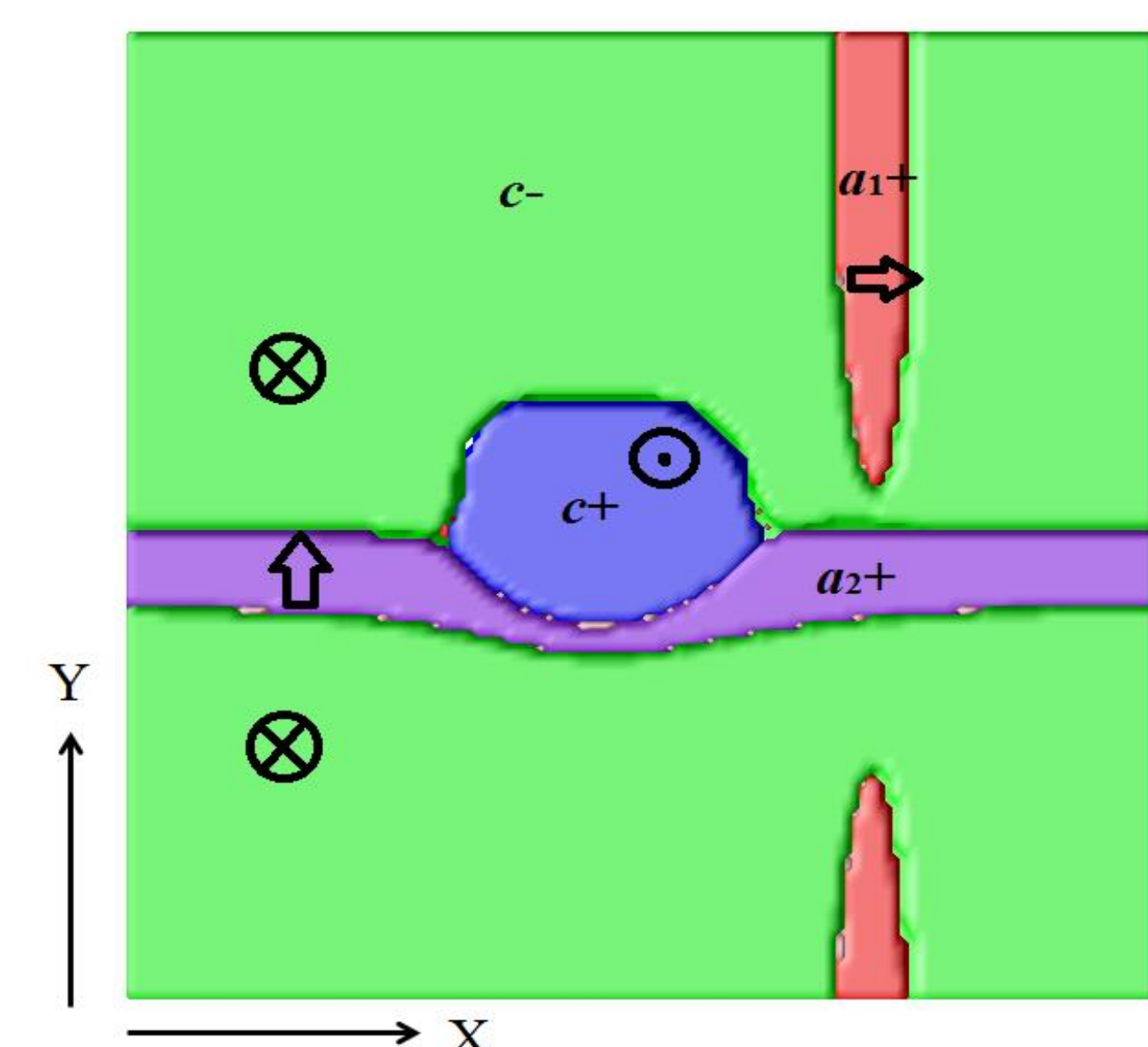
❖ $c+$ domain grows and interacts with the a_1+ domain, forming tail to tail configuration with charged domain wall.

❖ Part of the a_1+ domain switches to $c+$ domains, remaining a_1+ domains form triangularly shaped stripe a -domains to minimize the bound charges

❖ Growth of the switched c domain occurs through thinning and elongation of triangularly shaped stripe a -domains

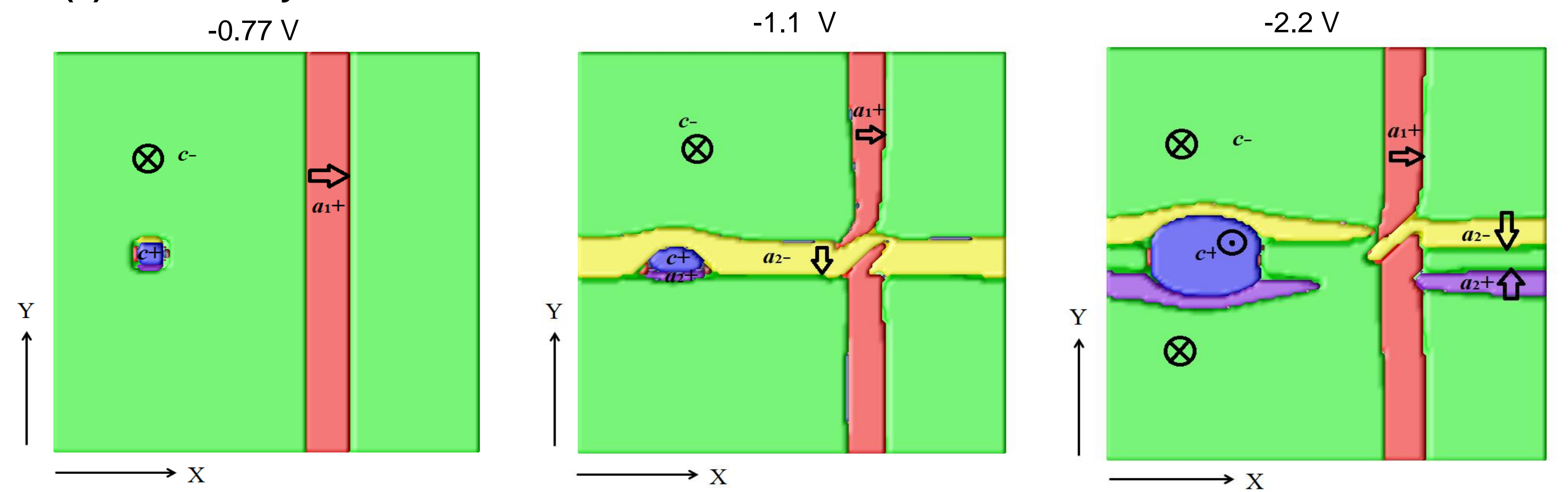
❖ Similar switching and growth process has been observed in experiments [5]

(4) Final domain structure with tip near domain wall



Switching under different strain and applied field

(1) Coherency strain -0.4%



❖ 3 different switching patterns are observed:

❖ Growth of c domains

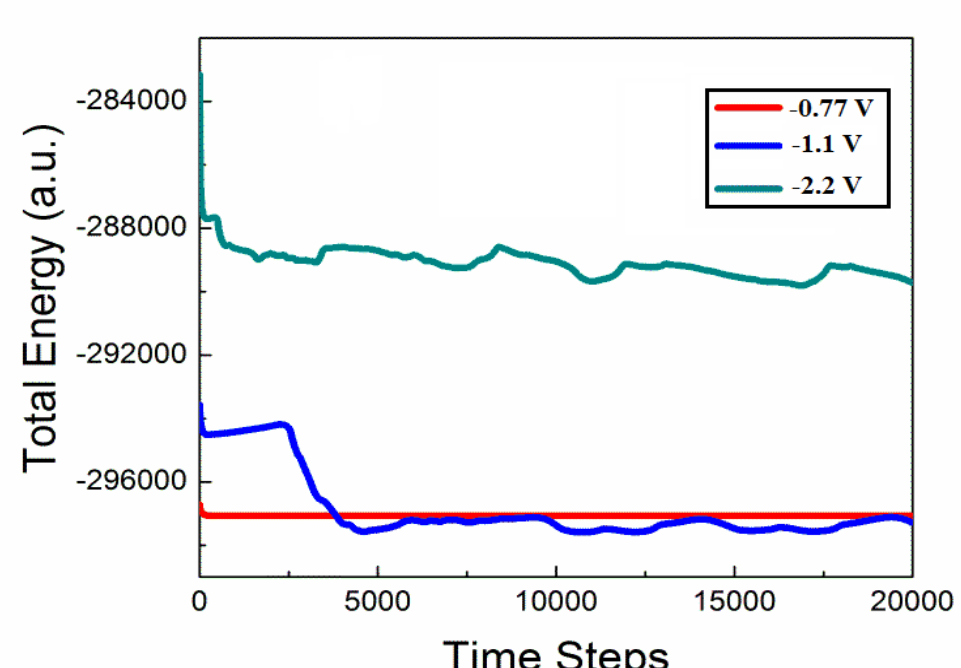
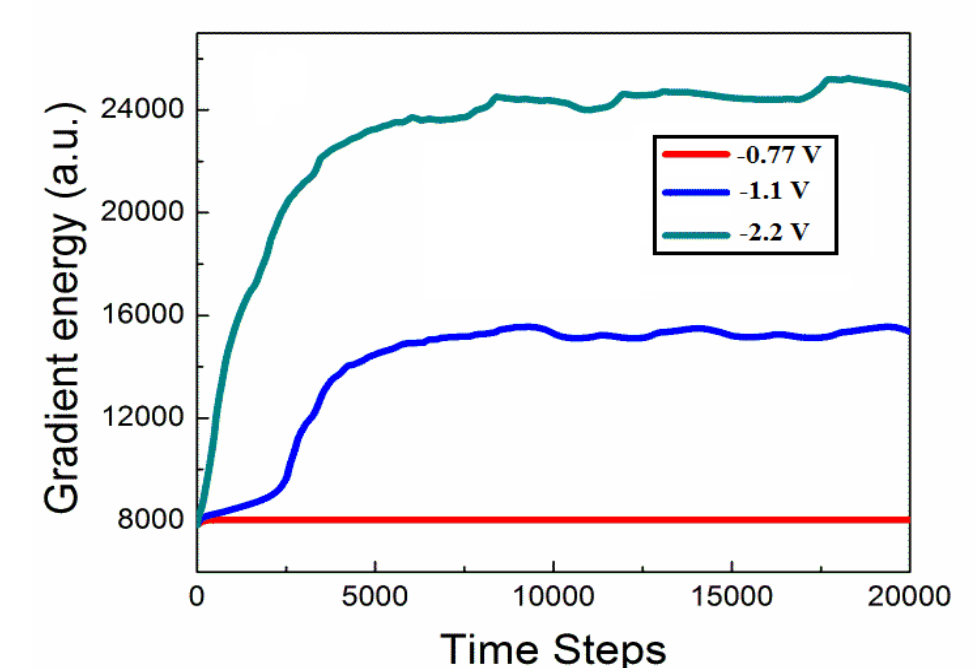
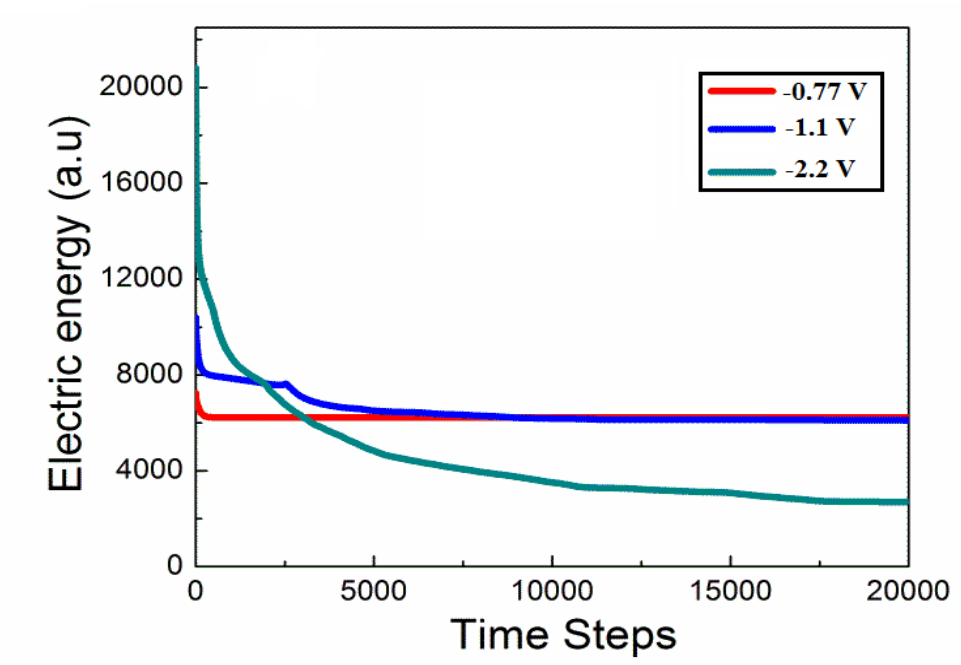
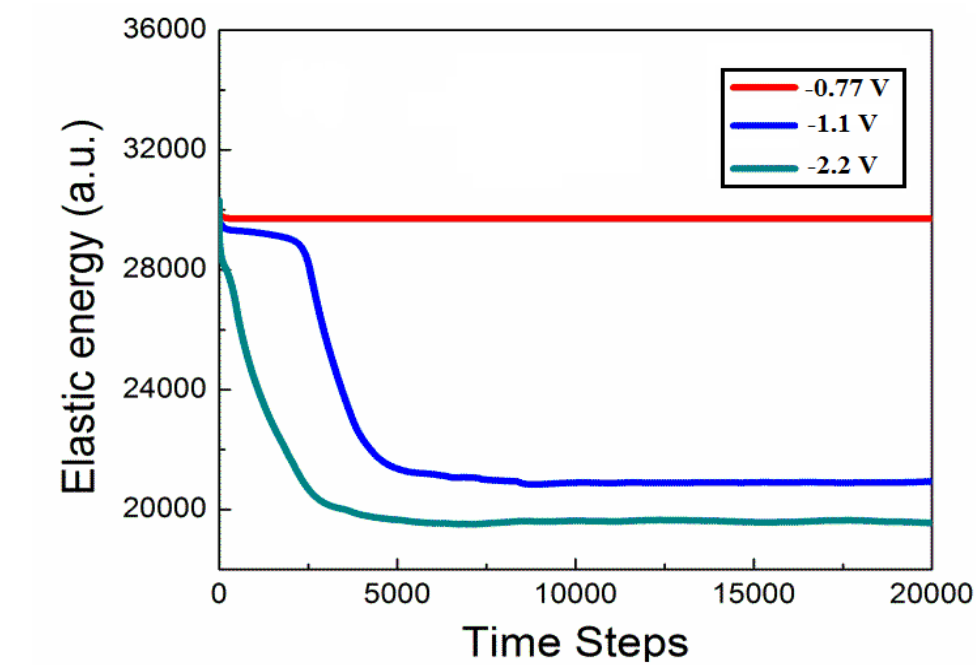
❖ Growth of c domain and one a domain

❖ Growth of c domain and two a domains

❖ Dominated by 180 degree switching

❖ Dominated by 90 degree switching

❖ Dominated by 180 degree switching again

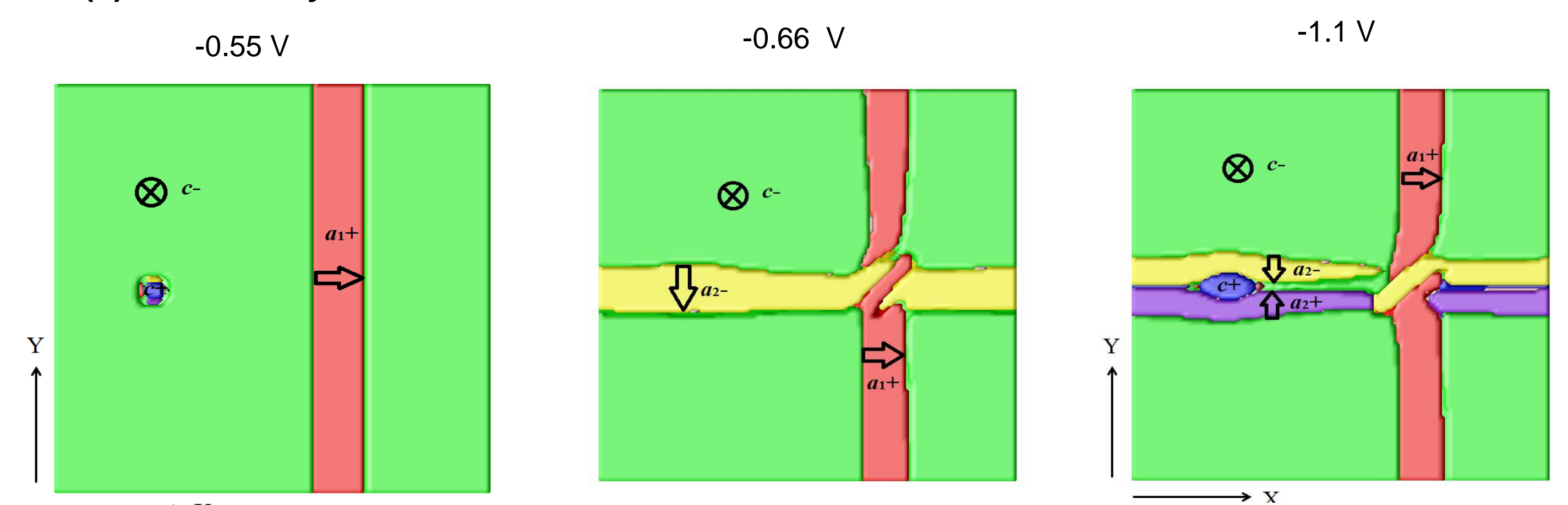


❖ With increasing bias, elastic energy drops with the growth of a domains

❖ Reductions in the elastic energy become smaller for the growth of additional a -domains

❖ The gradient energy increases dramatically as more and more domain walls are involved

(2) Coherency strain -0.3%



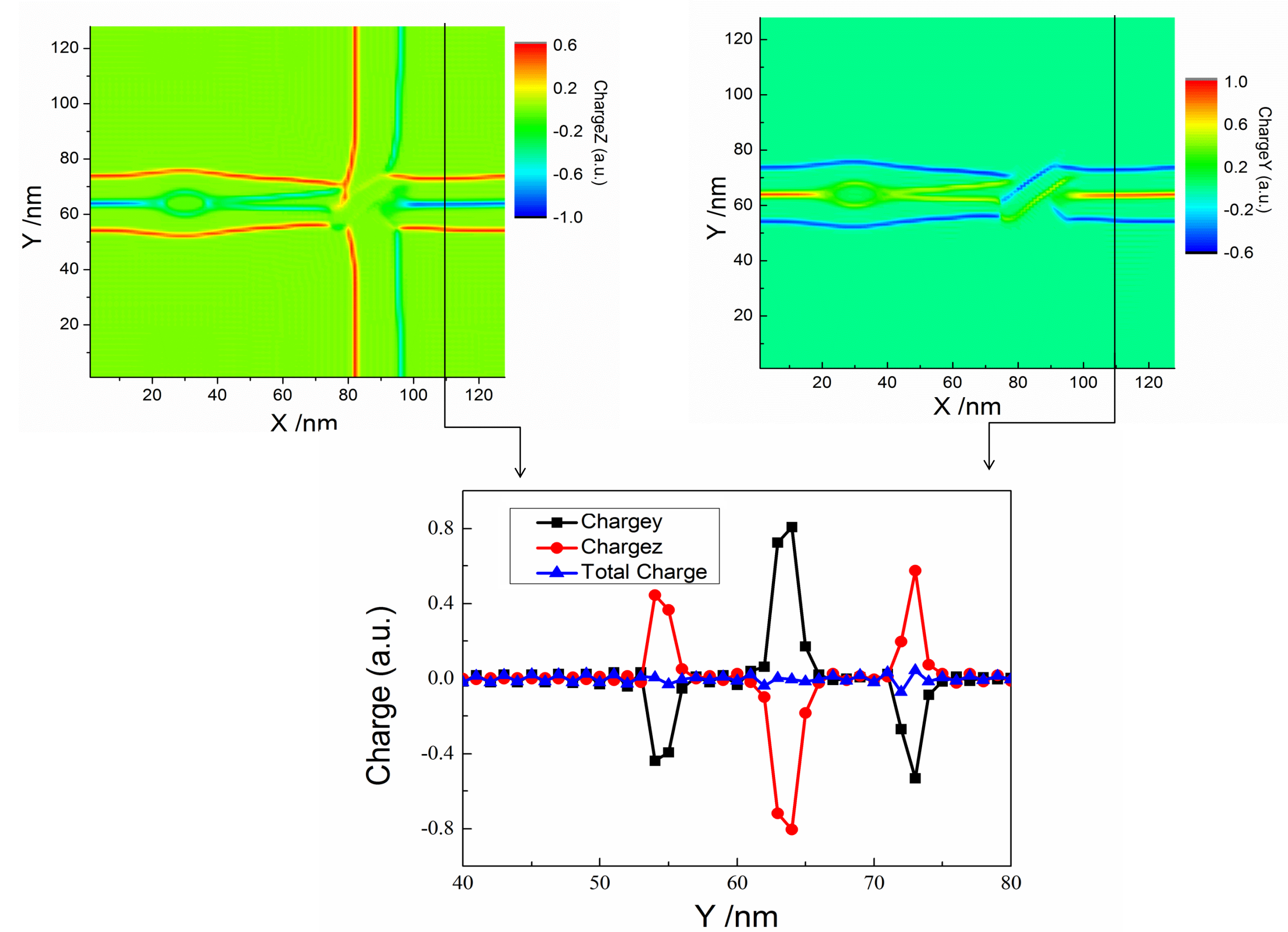
❖ The bias for 90 degree switching is decreased compare to -0.4% case

❖ When we increase the bias further, 90 degree switching from $c-$ to $a-$ is favored

❖ At intermediate bias, both 90 degree and 180 degree switching are favored

❖ 180 degree switching occur between the a domains

(3) Charge distribution with applying bias -1.1 V



❖ As the two a_2 -domains get close, the high P_2 gradient induces high charge

❖ This charge is compensated by chargez, which means the charge induced by P_3 gradient

❖ Thus 180° switching is favored between the two a_2 -domains due to charge compensation.

References

- (1) S. V. Kalinin, A. N. Morozovska, L. Q. Chen, and B. J. Rodriguez, *Rep. Prog. Phys.* **73** (5), 056502 (2010).
- (2) L. Chen, J. Ouyang, C. S. Ganpule, et al., *Appl. Phys. Lett.* **84** (2), 254 (2004).
- (3) Y. L. Li, S. Y. Hu, Z. K. Liu, et al. *Acta Mater.* **50** (2), 395 (2002).
- (4) L. Q. Chen and J. Shen, *Comput. Phys. Commun.*, **108** (2-3), 147 (1998).
- (5) J. K. Lee, G. Y. Shin, K. Song, et al. *Acta Mater.* **61** (18), 6765 (2013).
- (6) Z. J. Hong, J. Britson, J.-M. Hu and L.-Q. Chen, *Acta Mater.* **73** (2014): 75-82

Acknowledgements

This work is supported by the NSF Grant Nos. DMR-0820404, DMR-1006541, and DMR-1210588.

Conclusions

- ❖ 90° switching is favored during the local 180° switching
- ❖ Decreasing the magnitude of compressive strain could lead to an increase in elastic energy, which favors the growth of a -domains
- ❖ At relatively high bias, however, 180° switching is favored due to a rapid decrease in electrostatic energy.
- ❖ At medium bias, with low magnitude of compressive strain, due to charge compensation, 180° switching may occur far away from PFM tip near the two newly-grown a -domains.